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**Detection of dislocation by the moiré pattern in electron micrographs.** By HATSUJIRO HASHIMOTO, *Kyoto Technical University, Matsugasaki, Sakyo-ku, Kyoto, Japan* and RYOZI UYEDA, *Physical Institute, Nagoya University, Nagoya, Japan*

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By using an electron microscope of the highest resolution, Menter (1956) observed dislocations in direct images of 12 Å net spacings in crystal lattice of platinum and copper phthalocyanine. The present note deals with the detection of dislocations by the moiré pattern which can be observed by electron microscopes of lower resolution.

A moiré pattern is a coarse mesh pattern which is produced when two fine mesh are superimposed. The one-dimensional moiré pattern is easily demonstrated: draw two sets of equally spaced parallel lines on separate sheets of transparent paper, place one on the other and slowly rotate one with respect to the other; a changing moiré pattern is obtained. Moiré spacing  $s$  is given approximately by

$$s = d/\theta, \quad (1)$$

where  $d$  is the spacing of the parallel lines and  $\theta$  the angle of rotation. Moiré patterns are produced not only by rotating but also by inclining the meshes towards each other.

Moiré patterns on superposed atomic lattices were first reported by Mitsuishi, Nagasaki & Uyeda (1951), who observed moiré patterns on thin sheets of graphite crystal by an electron microscope.

Recently we found dislocations in moiré patterns on thin cupric sulphide crystals. Fig. 1(a) shows a region wherein the crystal lattice is almost perfect and Fig. 1(b) a region containing several dislocations. Figs. 1(c), (d) and (e) show highly magnified pictures of isolated and coupled dislocations. It is interesting that a single dislocation in a moiré pattern corresponds to a single dislocation in an atomic lattice. This fact is proved by the following demonstration: Fig. 2(a) is a model of net planes having a dislocation near its centre and in Fig. 2(b) there is no dislocation. Superpose these figures and slowly rotate one

with respect to the other. Then a single dislocation is observed in the moiré pattern independent of the angle of rotation, i.e. independent of the moiré spacing (Fig. 2(c) and (d)). Coupled dislocations (Fig. 1(d) and (e)) are also demonstrated by suitable models (Fig. 2(e) and (f)). In Fig. 3 many moiré dislocations of the same sign are appearing along a curve, indicating a small-angle tilt boundary (Burgers, 1939). It is most interesting that atomic dislocations which are too fine to be observed by present-day electron microscopes can be detected through the moiré pattern.

It was pointed out by Kainuma (1956) that if a moiré pattern is produced by two inclined crystal sheets, a screw dislocation nearly normal to the sheets produces a moiré dislocation also. By the selected-area diffraction method, however, it was proved that the present moiré pattern was caused by rotated crystal sheets, as was studied by Seki (1953) for a moiré pattern on sericite. The measured angle of rotation is  $\theta = 0.018$  rad. and the observed moiré spacing is  $s = 100$  Å. The net plane causing the moiré pattern is (11.0) of the cupric sulphide crystal and the spacing is  $d = 1.88$  Å. These values satisfy equation (1). Thus the moiré dislocations in Fig. 1 are without doubt caused by edge dislocations nearly normal to the crystal sheets.

#### References

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**Unit-cell parameters and space group of bisacetylacetoneberyllium.** By G. J. BULLEN, *Department of Chemistry, University College, Achimota, Gold Coast, British West Africa*

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Bisacetylacetoneberyllium,  $\text{Be}(\text{C}_5\text{H}_7\text{O}_2)_2$ , was reported by Jaeger (1914) to be monoclinic prismatic with

$$a:b:c = 1.4765:1:1.3592; \quad \beta = 52^\circ 20'.$$

From Weissenberg photographs taken using  $\text{Cu K}\alpha$  radiation a different unit cell has now been chosen as preferable. This cell has the dimensions

$$a = 13.45 \pm 0.06, \quad b = 11.30 \pm 0.04, \quad c = 7.74 \pm 0.05 \text{ Å}; \\ \beta = 100.8 \pm 0.1^\circ,$$

$a$  and  $c$  being the two shortest distances between lattice points in the (010) plane. Jaeger's  $c$  is double the new  $c$  and his  $a$  is the new  $[\bar{1}01]$ . Calculating from the new cell dimensions, we get

$$[\bar{1}01]:b:2c = 1.48:1:1.37$$

and the angle between  $[\bar{1}01]$  and  $c$  is  $52.2^\circ$ , agreeing well with Jaeger's determination.

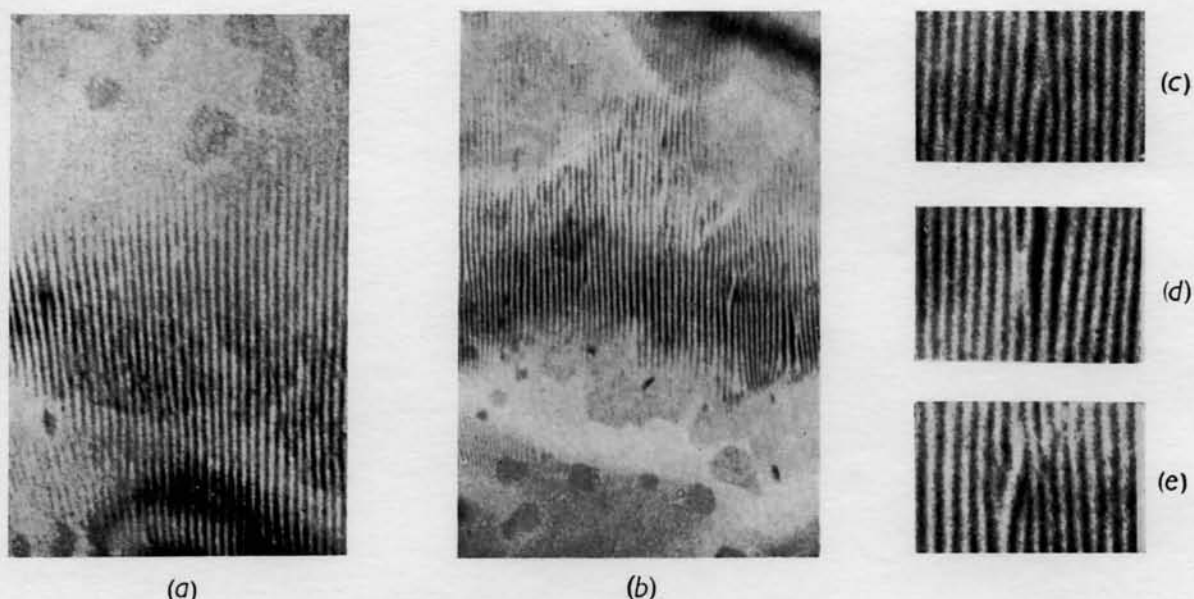


Fig. 1. Electron micrographs of moiré patterns on cupric sulphide crystals.

- (a) Nearly perfect region.  $\times 120,000$ .  
 (b) Region with several dislocations.  $\times 85,000$ .  
 (c), (d), (e). Single and coupled dislocations.  $\times 220,000$ .

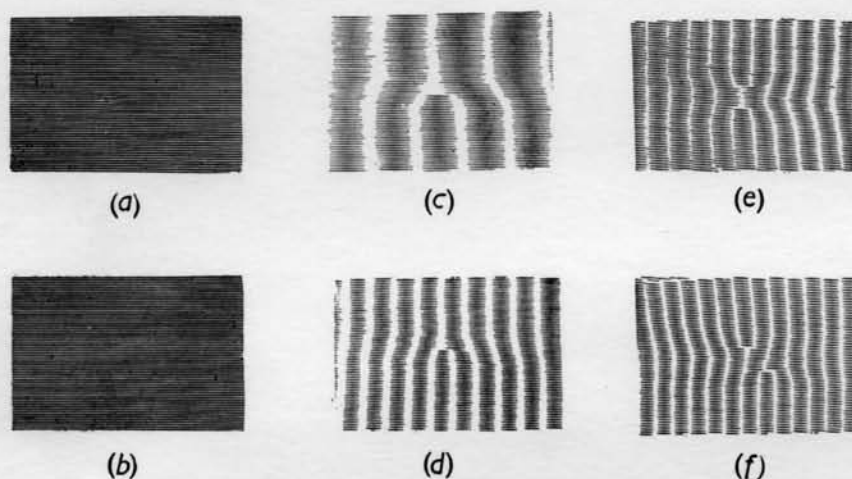


Fig. 2. Model of moiré dislocation.

- (a) and (b) correspond to net planes with and without dislocation, respectively.  
 (c) and (d) are superposition of (a) and (b) with angles of rotation of  $3.2^\circ$  and  $6.2^\circ$ , respectively.  
 (e) and (f) are models of coupled moiré dislocations.



Fig. 3. Moiré pattern on tilt boundary.  $\times 100,000$ .